

Molecular Dynamics

Goal: **trajectories** $x(t)$ that sample the different conformations of the system so that we can calculate correct **thermodynamic observables**.

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Molecular Dynamics Algorithm

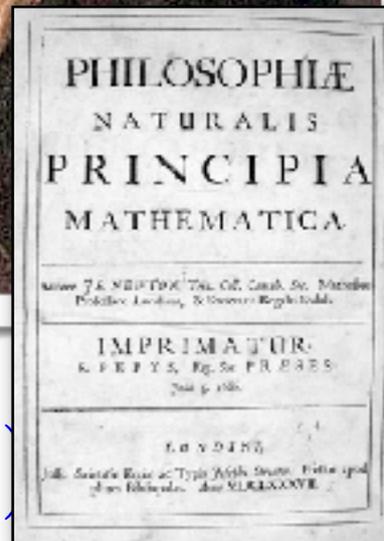
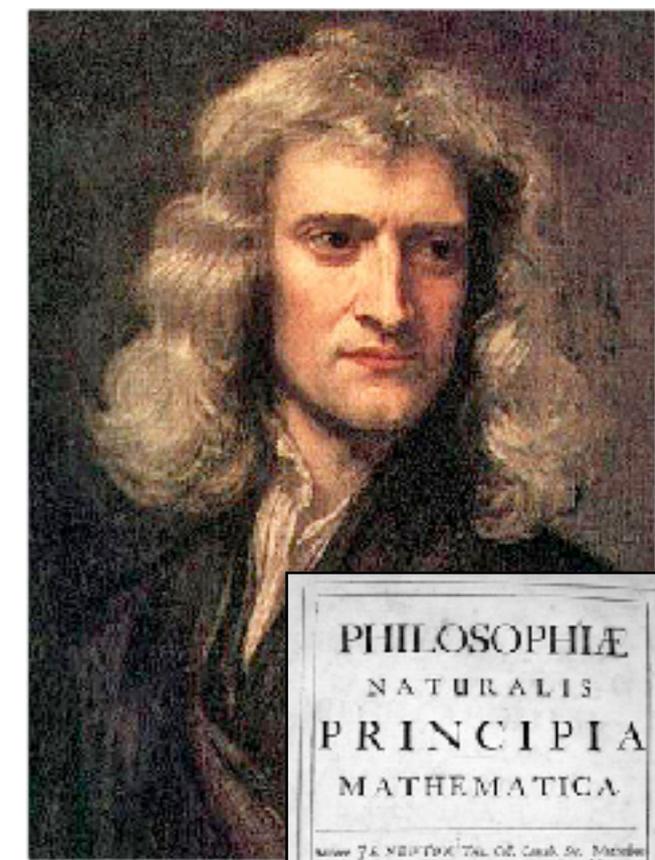
- equations of motion for N particles

$$\mathbf{F}_i = m_i \mathbf{a}_i = m_i \frac{d^2 \mathbf{r}}{dt^2} = m_i \ddot{\mathbf{r}}_i$$

- potential energy

$$U(\mathbf{r}_1, \dots, \mathbf{r}_N) = U_{\text{bonded}}(\{\mathbf{r}_i\}) + U_{\text{nonbonded}}(\{\mathbf{r}_i\})$$

source: Wikipedia



$$U_{\text{bond}}(\mathbf{r}_j, \mathbf{r}_k) = \frac{K}{2} (|\mathbf{r}_j - \mathbf{r}_k| - b_0)^2$$

$$U_{\text{angles}}(\mathbf{r}_i, \mathbf{r}_j, \mathbf{r}_k) \quad U_{\text{dih}}(\mathbf{r}_i, \mathbf{r}_j, \mathbf{r}_k, \mathbf{r}_l)$$

$$U_{\text{elec}}(\mathbf{r}_j, \mathbf{r}_k) = \frac{1}{4\pi\epsilon_0} \frac{q_j q_k}{|\mathbf{r}_j - \mathbf{r}_k|}$$

$$U_{\text{vdW}}(\mathbf{r}_j, \mathbf{r}_k) = 4\epsilon \left[\left(\frac{\sigma}{r_{jk}} \right)^{12} - \left(\frac{\sigma}{r_{jk}} \right)^6 \right]$$

- forces

$$\mathbf{F}_i = -\nabla_i U(\mathbf{r}_1, \dots, \mathbf{r}_i, \dots, \mathbf{r}_N)$$

Basic MD program

initialize()

$t=0$

while $t < t_{\max}$

$\mathbf{f} = \text{forces}(\mathbf{r})$

$\mathbf{r} = \text{integrate}(\mathbf{r}, \mathbf{f})$

write_trajectory(\mathbf{r})

$t = t + dt$

Basic workflow

1. prepare starting conformation
2. run equilibration
3. run production
 - compute estimators for observables $a(\mathbf{r}, \mathbf{v})$
 - save trajectory (\mathbf{r} , maybe \mathbf{v} and \mathbf{F})
4. analyse trajectory/output files

Force calculation

initialize()

$t=0$

while $t < t_{\max}$

$\mathbf{f} = \text{forces}(\mathbf{r})$

$\mathbf{r} = \text{integrate}(\mathbf{r}, \mathbf{f})$

***write_trajectory*(\mathbf{r})**

$t = t + dt$

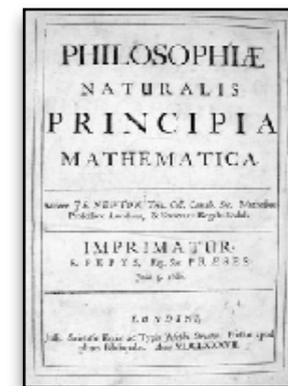
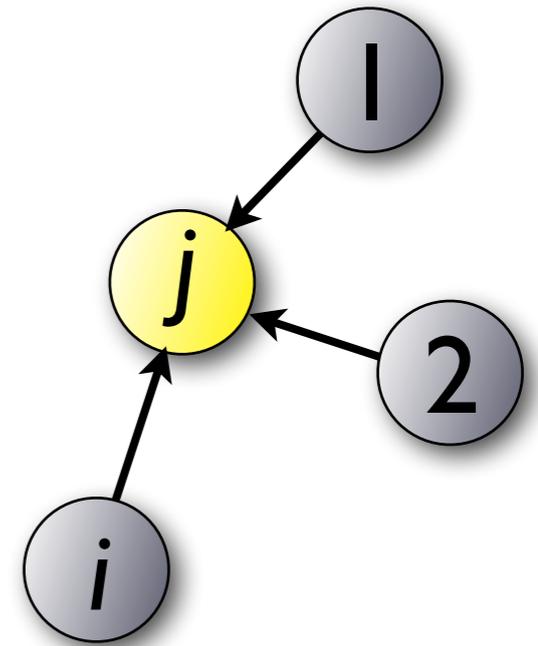
Force calculation

$$\mathbf{F}_i = -\nabla_i U(\mathbf{r}_1, \dots, \mathbf{r}_i, \dots, \mathbf{r}_N)$$

For pair-potentials $v(|\mathbf{r}_j - \mathbf{r}_i|) = v(r_{ij})$

$$\mathbf{F}_j = \sum_i \underbrace{[-\nabla_i v(r_{ij})]}_{\mathbf{f}_{ij}} = - \sum_i \frac{\mathbf{r}_{ij}}{r_{ij}} \frac{\partial}{\partial r_{ij}} v(r_{ij})$$

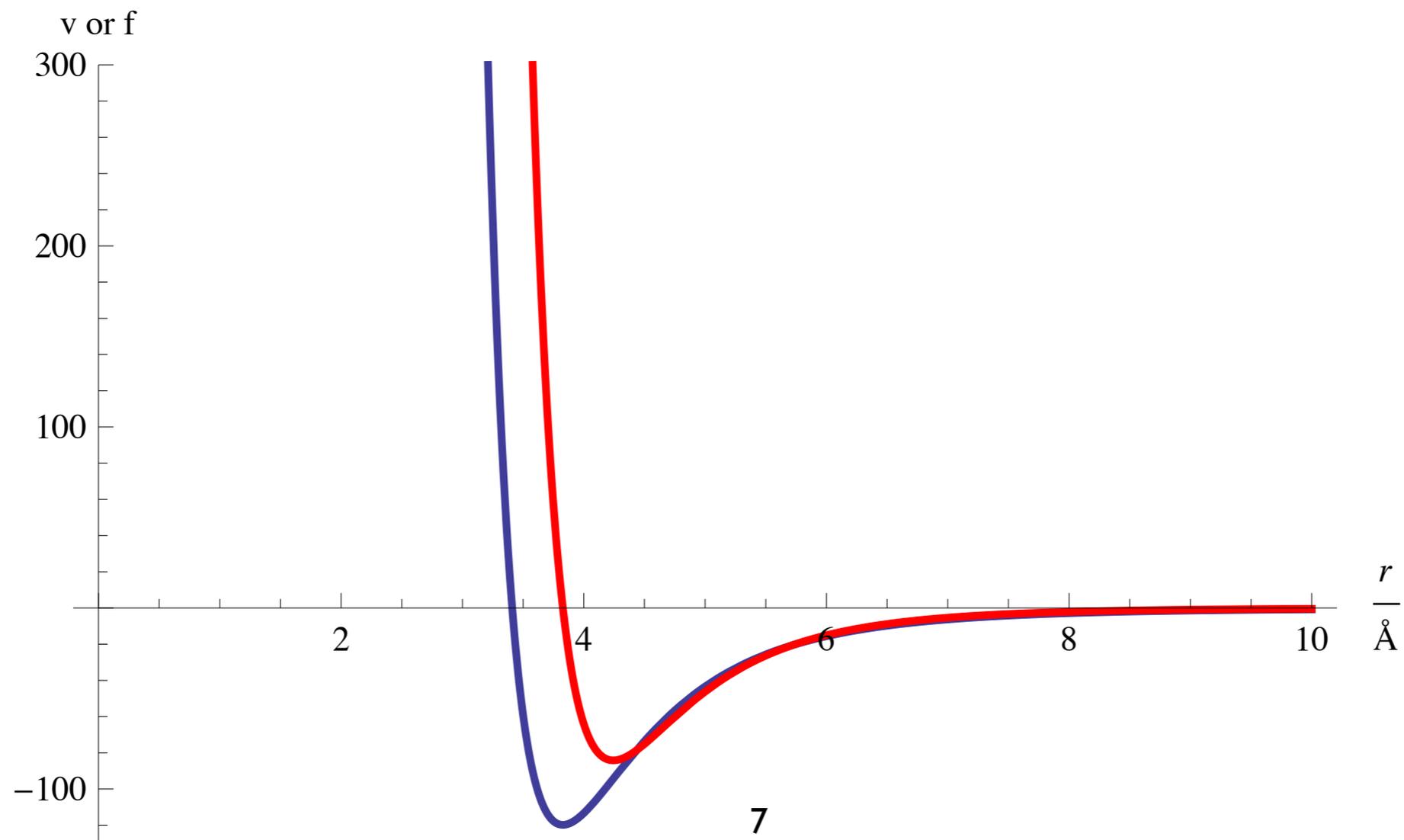
$$\mathbf{f}_{ij} = -\mathbf{f}_{ji} \quad (\text{Newton's 3rd law})$$



Lennard-Jones Potential

$$v(r_{ij}) = 4\epsilon \left[\left(\frac{\sigma}{r_{ij}} \right)^{12} - \left(\frac{\sigma}{r_{ij}} \right)^6 \right]$$

$$\mathbf{f}(r_{ij}) = -24\epsilon \frac{\mathbf{r}_{ij}}{r_{ij}^2} \left[2 \left(\frac{\sigma}{r_{ij}} \right)^{12} - \left(\frac{\sigma}{r_{ij}} \right)^6 \right]$$



Integrating the Equations of Motion

initialize()

$t=0$

while $t < t_{\max}$

$\mathbf{f} = \text{forces}(\mathbf{r})$

$\mathbf{r} = \text{integrate}(\mathbf{r}, \mathbf{f})$

write_trajectory(\mathbf{r})

$t = t + dt$

Integrating the Equations of Motion

$$m_i \ddot{\mathbf{r}}_i = \mathbf{F}_i$$

- have forces \mathbf{F}_i and current positions $\mathbf{r}_i(t)$
- want: positions $\mathbf{r}_i(t+\Delta t)$

Velocity Verlet Integrator

$$\mathbf{v}(t + \frac{\Delta t}{2}) = \mathbf{v}(t) + \frac{\Delta t}{2} \frac{\mathbf{F}(t)}{m} \quad (1)$$

$$\mathbf{r}(t + \Delta t) = \mathbf{r}(t) + \Delta t \mathbf{v}(t + \frac{\Delta t}{2}) \quad (2)$$

$$\mathbf{v}(t + \Delta t) = \mathbf{v}(t + \frac{\Delta t}{2}) + \frac{\Delta t}{2} \frac{\mathbf{F}(t + \Delta t)}{m} \quad (3)$$

- gives both positions and velocities
- one force evaluation per time step
- see Module 10 ODEs

Good integrators

- speed? – not very relevant
- accuracy for large time steps Δt
- **energy conservation** (for $\partial H/\partial t = 0$)
 - short term
 - long term (more important)
- time reversible (Newton's EOMs!)
- phase-space area preserving (Hamiltonian dynamics!)
- accurately predict the trajectories of all particles for short and long times?

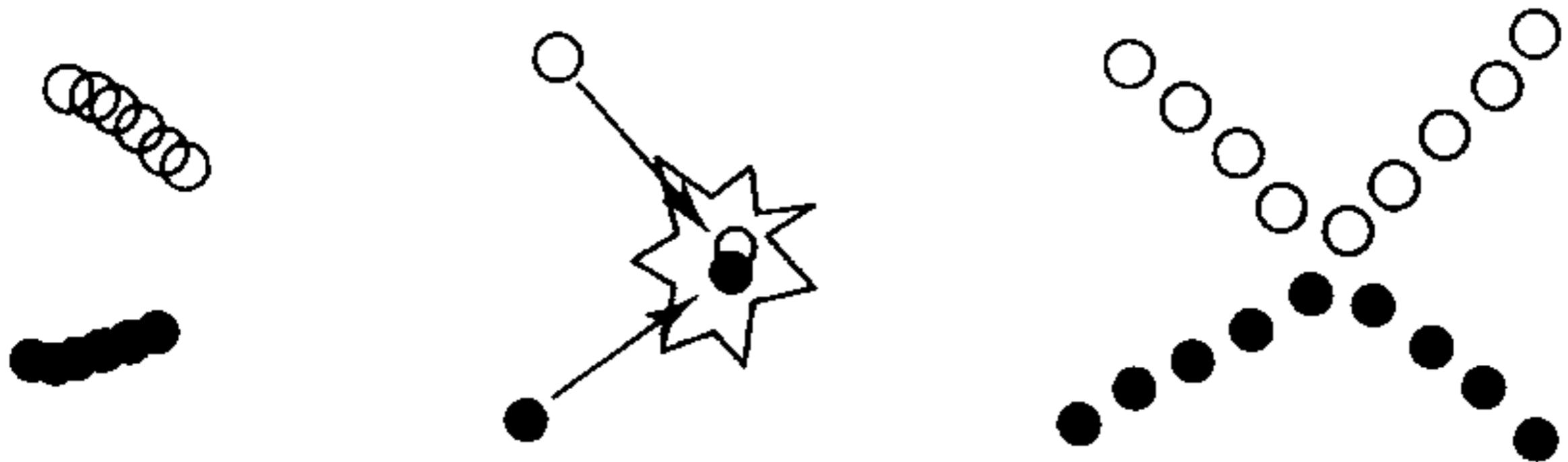
Good integrators

- speed/memory? – not very relevant
- accuracy for large time steps Δt
- energy conservation (for $\partial H/\partial t = 0$)
 - short term
 - long term (more important)
- time reversible (Newton's EOMs!)
- phase-space area preserving (Hamiltonian dynamics!)
- accurately predict the trajectories of all particles for short and long times?

Verlet

- fast, small memory requirements
- not very high accuracy for long time steps
- fair short- and good long term energy conservation
- time reversible and area preserving (“symplectic”)

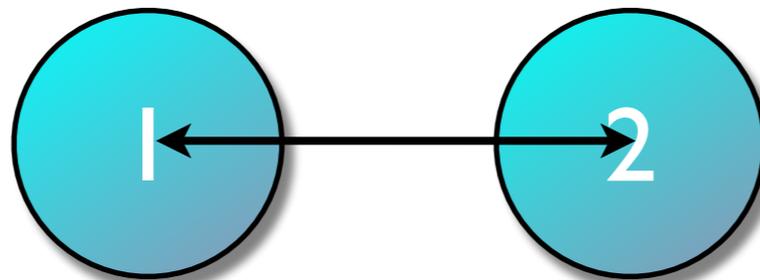
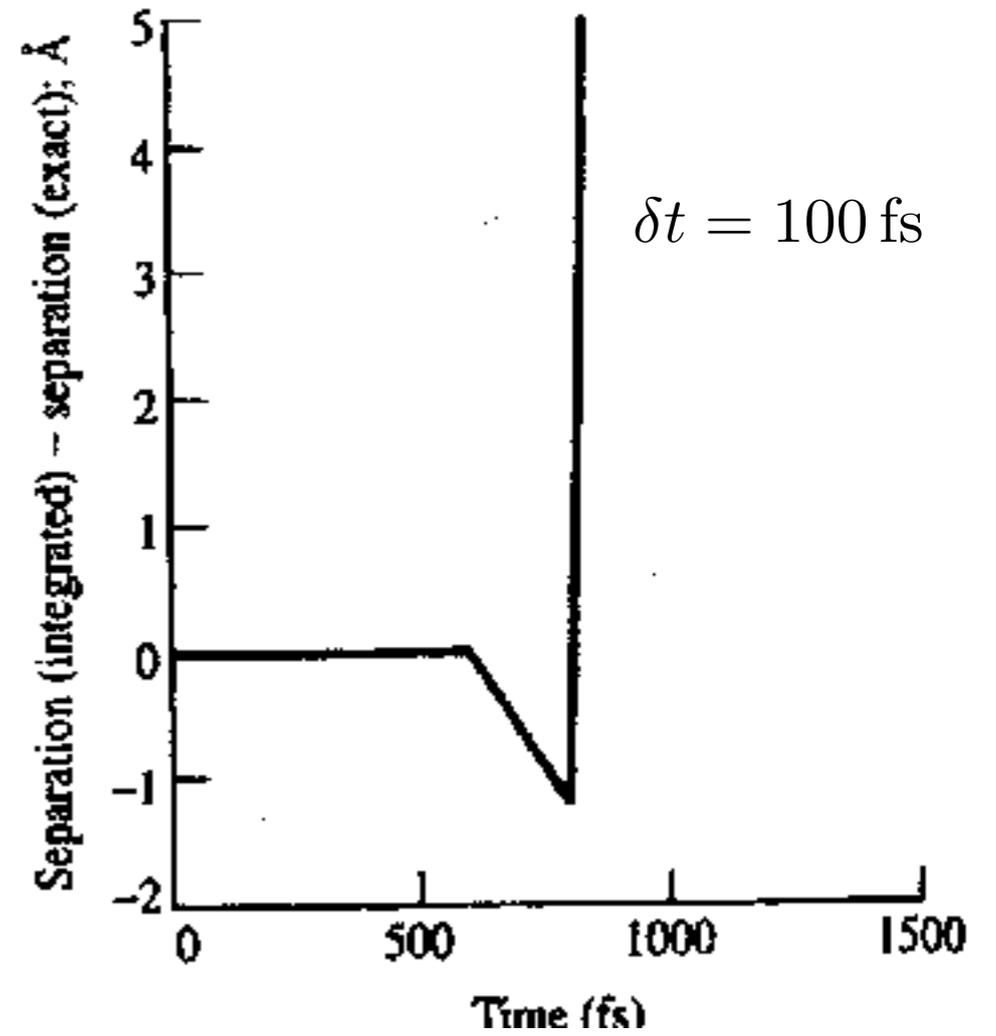
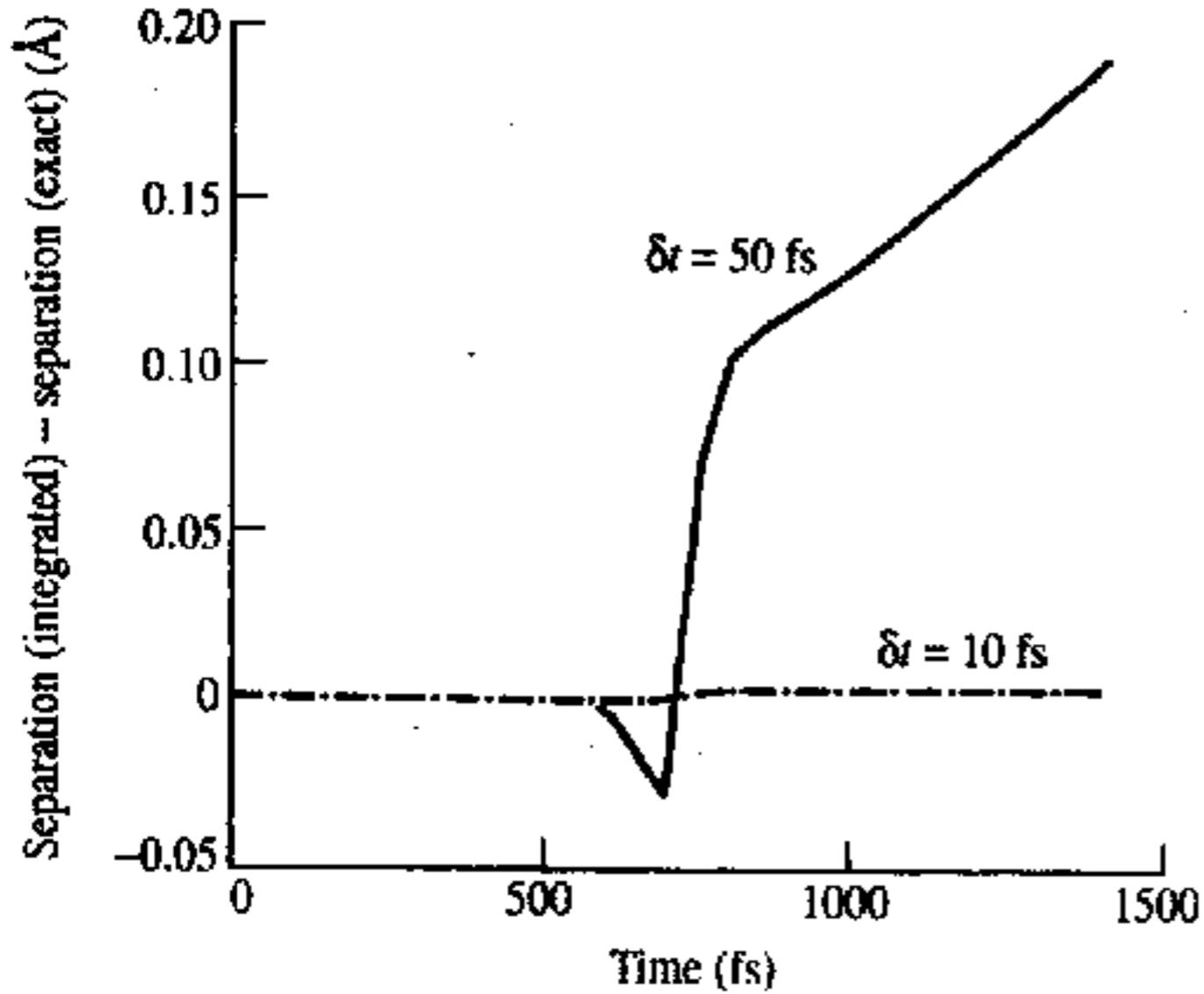
Time step



from Leach, 2001

- *rule of thumb for Verlet-type integrators: 5 steps per period (typically 1-2 fs in biomolecular systems)*

Time step



Thermodynamic observables

- **Macroscopic thermodynamic quantities** (temperature T , pressure P , volume V , particle number N , heat capacities, dipole moments, magnetization, ...) can be written as **averages** over **functions that depend on microscopic positions and velocities** (“estimator”).

$$A = \langle a \rangle$$

$$A = \langle a(\mathbf{x}_1, \mathbf{x}_2, \dots, \mathbf{x}_N, \mathbf{v}_1, \mathbf{v}_2, \dots, \mathbf{v}_N) \rangle$$

$$A = \langle a(\mathbf{x}_1(t), \mathbf{x}_2(t), \dots, \mathbf{x}_N(t), \mathbf{v}_1(t), \mathbf{v}_2(t), \dots, \mathbf{v}_N(t)) \rangle_t$$

$$A = \frac{1}{N_t} \sum_{i=1}^{N_t} a(\mathbf{x}_1(t), \mathbf{x}_2(t), \dots, \mathbf{x}_N(t), \mathbf{v}_1(t), \mathbf{v}_2(t), \dots, \mathbf{v}_N(t))$$

Energy E

$$E = \langle \mathcal{H} \rangle$$

$$\mathcal{H} = T_{\text{kin}}(t) + U(t)$$

$$= \sum_{i=1}^N \frac{1}{2} m_i \mathbf{v}_i^2 + \sum_{i < j}^{N-1} U_{\text{LJ}}(r_{ij})$$

$$r_{ij} := |\mathbf{r}_j - \mathbf{r}_i|$$

\mathcal{H} : “Hamiltonian” (total instantaneous energy)

Energy conservation: $\frac{d\mathcal{H}}{dt} = 0$

Temperature T

$$T_{\text{kin}}(t) = \frac{1}{2} \sum_{i=1}^N m_i \mathbf{v}_i(t)^2 \quad \langle T_{\text{kin}} \rangle = N \frac{3}{2} k_B T$$

1/2 kT per degree of freedom

$$T = \frac{2 \langle T_{\text{kin}} \rangle}{3 k_B N}$$

Boltzmann's constant

$$k_B = 1.38064852 \times 10^{-23} \text{ J} \cdot \text{K}^{-1}$$

$$\mathcal{T}(t) = \frac{2T_{\text{kin}}(t)}{3k_B N}$$

$$T = \langle \mathcal{T}(t) \rangle_t$$

$$A = \frac{1}{N_t} \sum_{i=1}^{N_t} a(\mathbf{x}_1(t), \mathbf{x}_2(t), \dots, \mathbf{x}_N(t), \mathbf{v}_1(t), \mathbf{v}_2(t), \dots, \mathbf{v}_N(t))$$

Pressure P

- pressure can be derived from a version of the *virial theorem* (... another time)

$$PV = Nk_B T + \frac{w}{3}$$

“virial” $w := \left\langle \sum_{i < j}^{N-1} \mathbf{r}_{ij} \cdot \mathbf{f}_{ij} \right\rangle$ $\mathbf{F}_i = \sum_{j \neq i} \mathbf{f}_{ij}$

$$\mathcal{P}(t) = V^{-1} \left[Nk_B \mathcal{T}(t) + \frac{1}{3} \sum_{i < j}^{N-1} \mathbf{r}_{ij}(t) \cdot \mathbf{f}_{ij}(t) \right]$$

$$P = \langle \mathcal{P}(t) \rangle_t$$

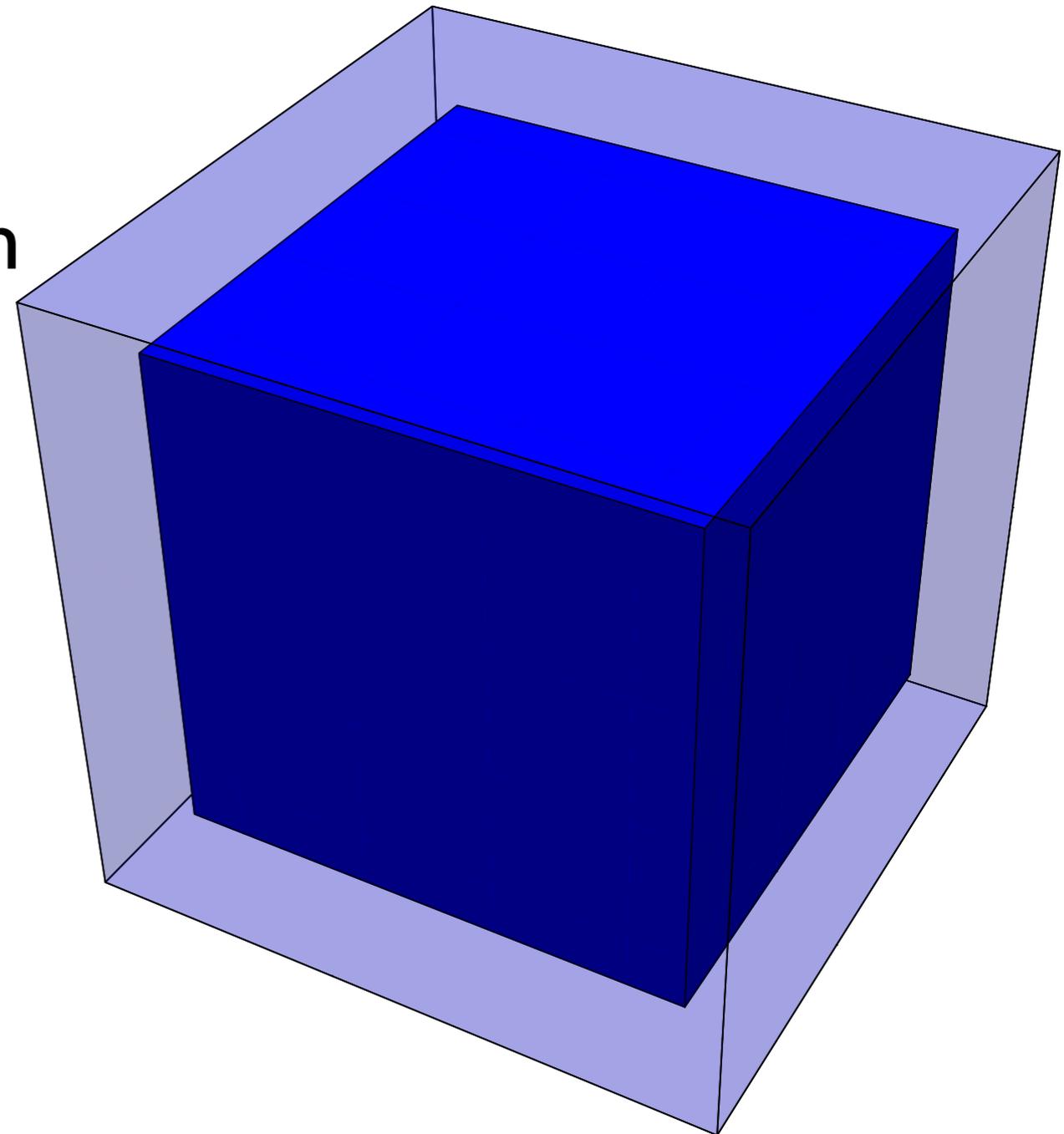
Summary

- outline of the MD algorithm
- use of periodic boundary conditions
- potential truncation
- integrators
- calculating macroscopic observables from microscopic estimators

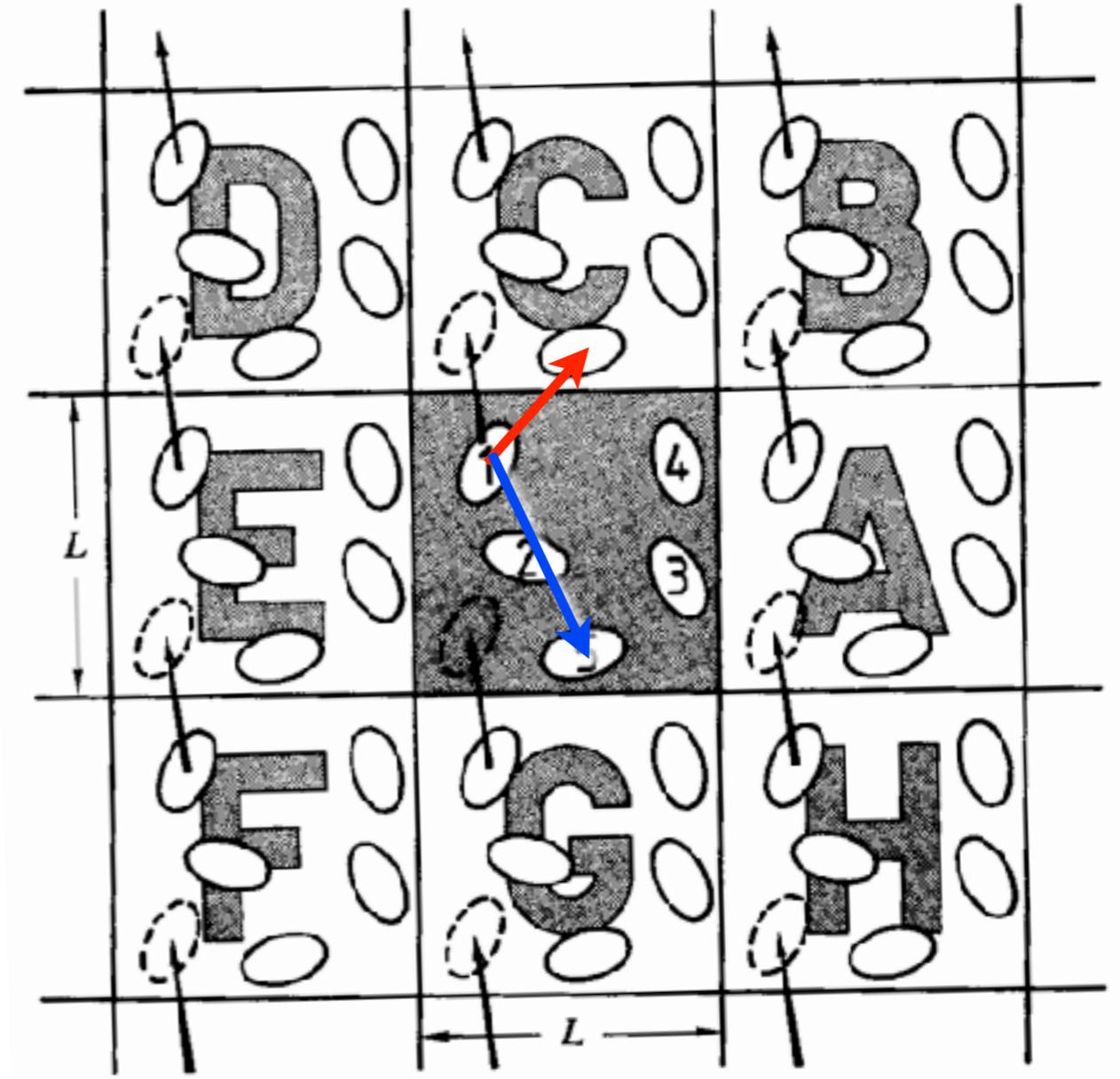
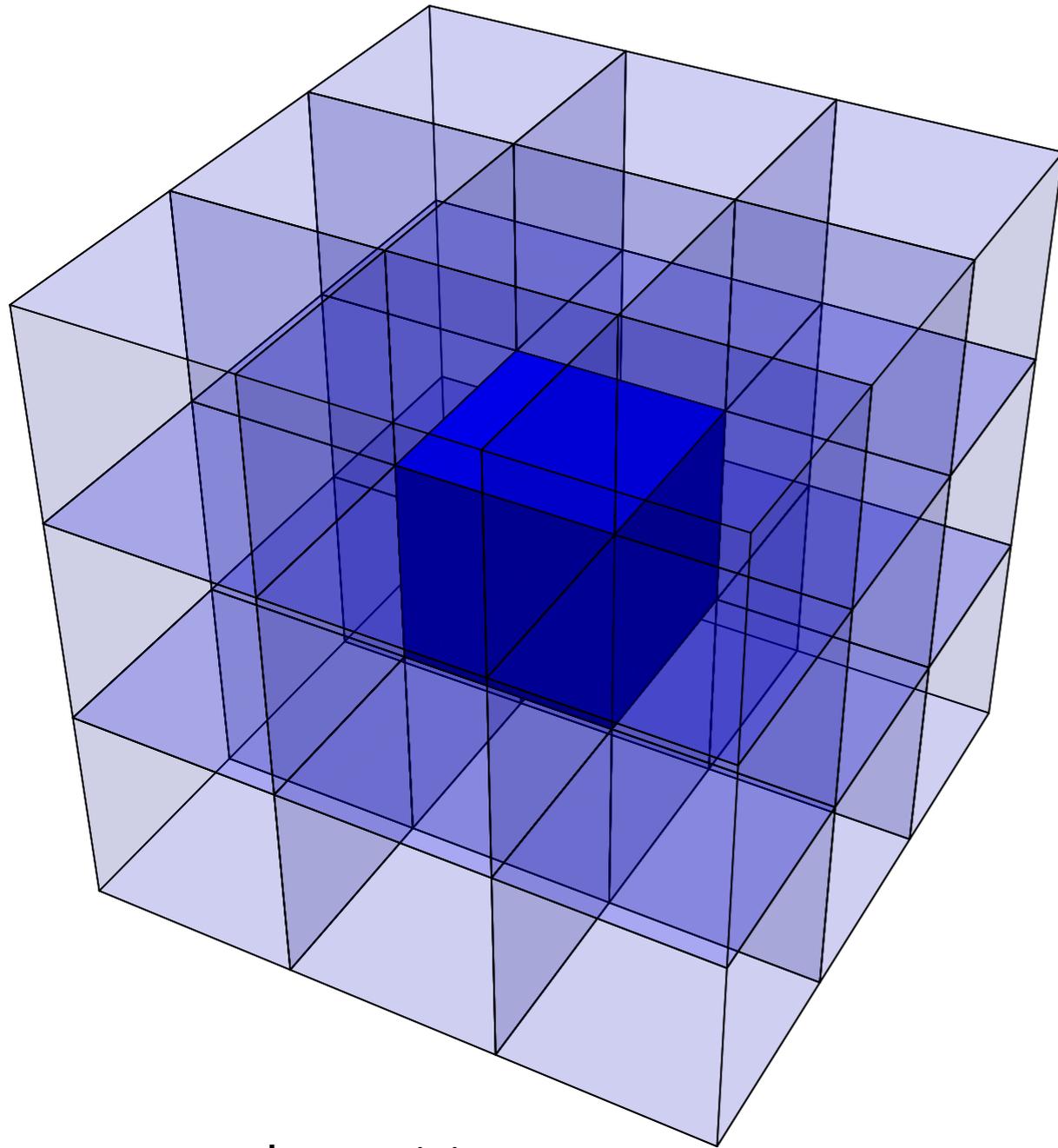
Appendix

Small systems: Surface effects

- In a small system, “most” particles are “near: the surface.
- “near” = typical interaction distance $d \sim 3\sigma$
- $N_{\text{surf}}/N \sim N^{-1/3}$
- *Problem* when we are interested in **bulk properties**.



Small systems: Periodic boundaries



box centered on origin:

$$\mathbf{r}_i \leftarrow \mathbf{r}_i - L \text{rint} \left(\frac{\mathbf{r}_i}{L} \right)$$

$$\mathbf{r}_{ij} \leftarrow \mathbf{r}_{ij} - L \text{rint} \left(\frac{\mathbf{r}_{ij}}{L} \right)$$

from Allen & Tildesley, 1987

Periodic boundaries: Potential Issues

- simulate *infinite* system so must handle an infinite amount of interactions (though often truncation is permissible)
- spurious correlations/ordering
- only fluctuations allowed with lattice periodicity, and max wavelength $\lambda=L$
 - ➔ phase transitions with long wavelength fluctuations problematic
- non-isotropic pair distribution function

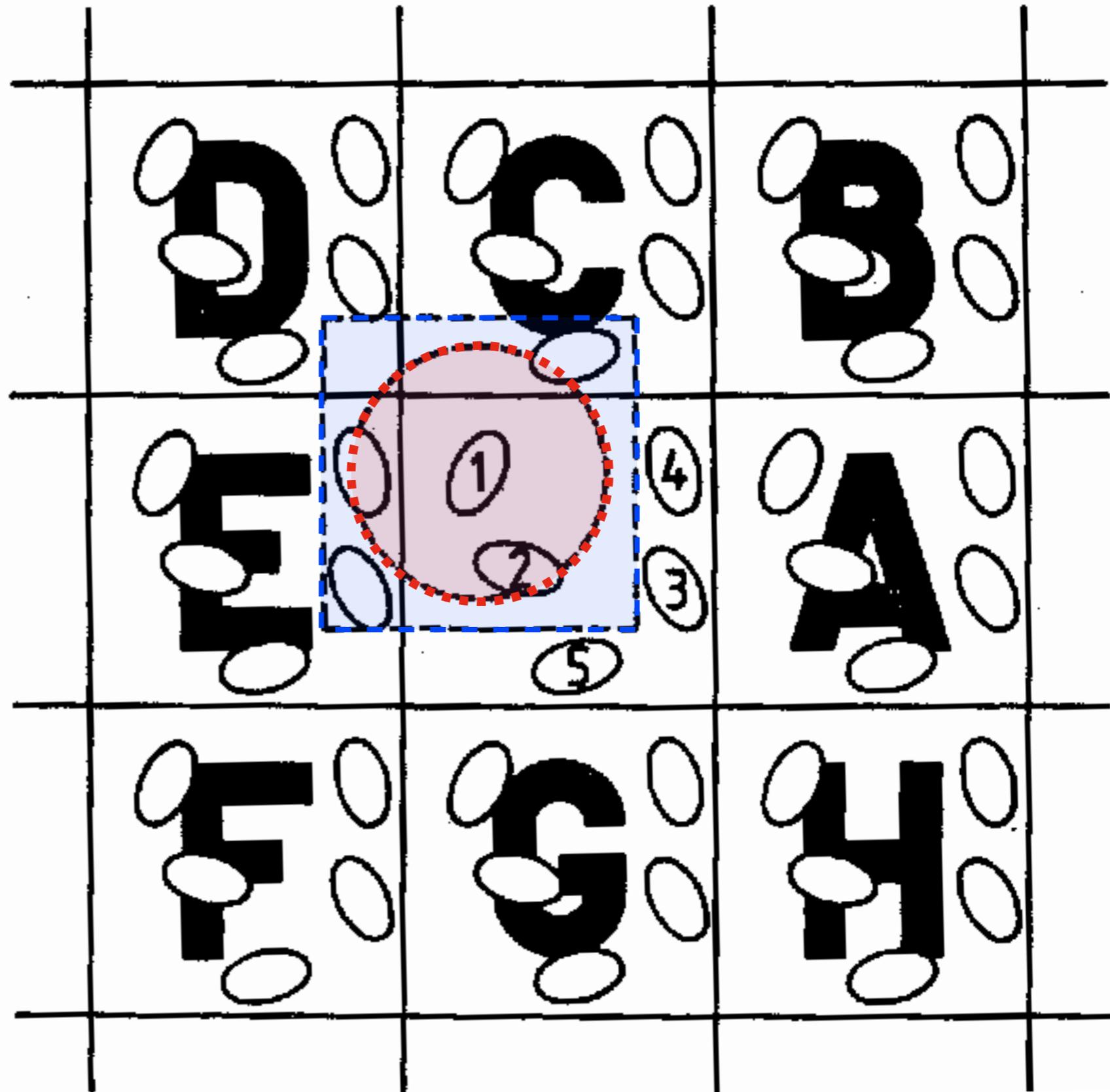
... but generally works rather well!

Periodic boundaries: Interactions

- pair wise $v(r_{ij}) \sim N^2$
- minimum image convention
- truncation at **cutoff R_c**

$v(r)$ must fall off faster than r^{-3}

$$R_c < L/2$$

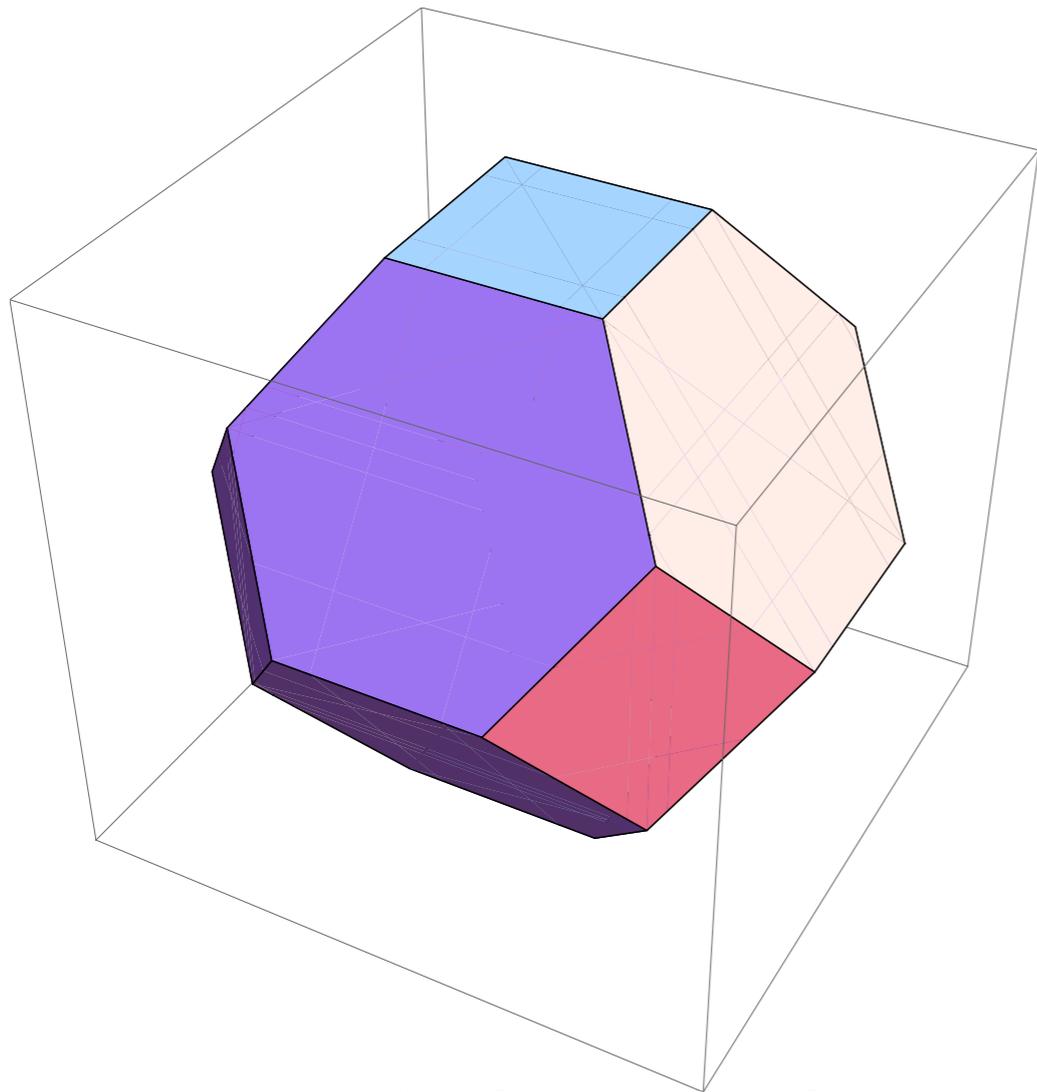


Truncation of potentials

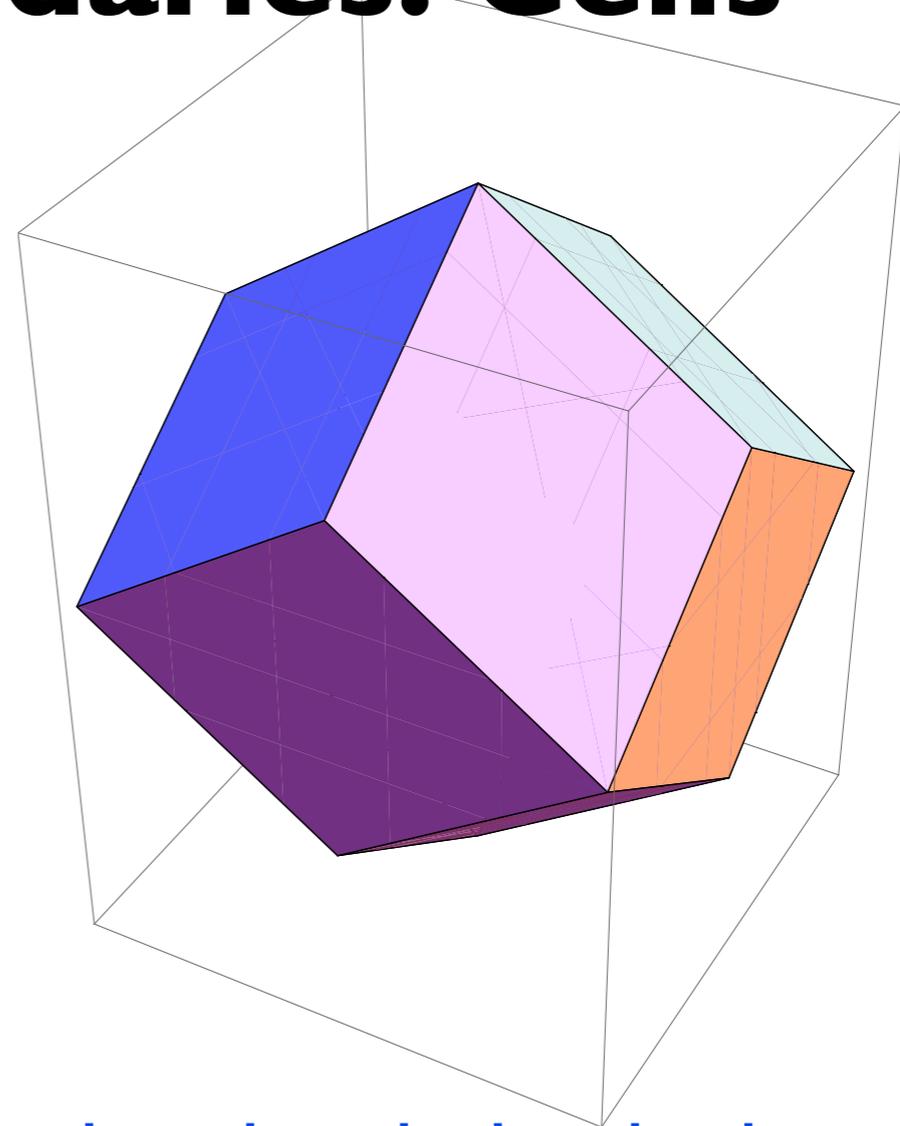
(proper periodic treatment: advanced topic)

- simple truncation $u^{\text{trunc}}(r) = \begin{cases} u(r) & r \leq R_c \\ 0 & r > R_c \end{cases}$
- truncation and shift $u^{\text{shift}}(r) = \begin{cases} u(r) - u(R_c) & r \leq R_c \\ 0 & r > R_c \end{cases}$
- minimum image convention

Periodic boundaries: Cells

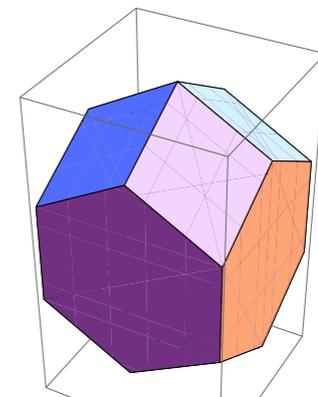
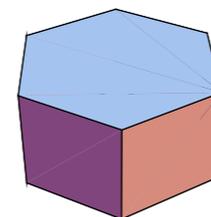
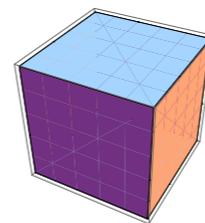


truncated octahedron
($0.77 d^3$)



rhombic dodecahedron
($0.71 d^3$)

- + cube/parallel epiped (d^3)
- + hexagonal prism
- + elongated rhombic dodecahedron



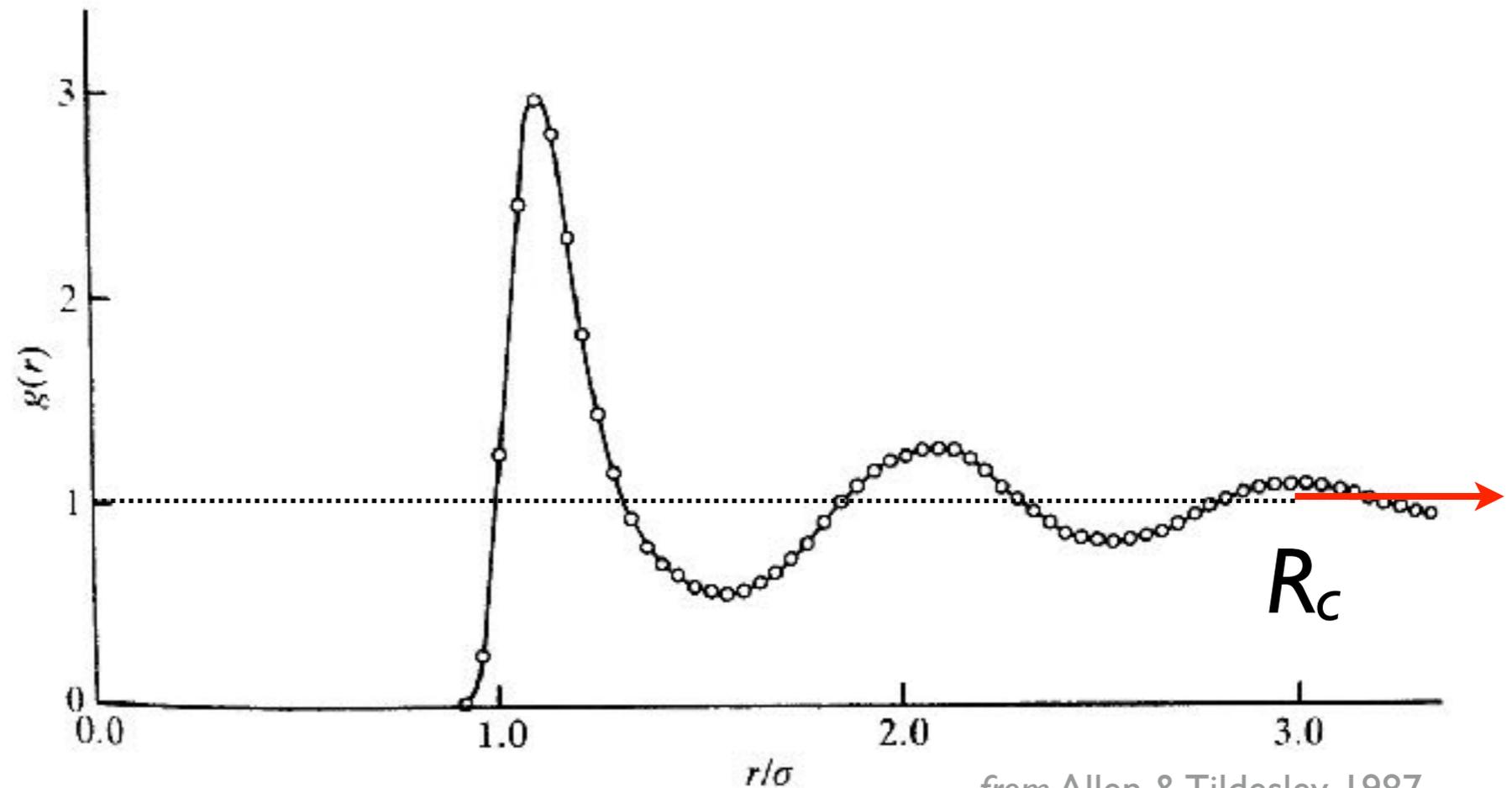
Periodic boundaries: Tail correction

$$U_{\text{tot}} = \frac{1}{2} N \rho \int_0^{\infty} u(r) g(r) 4\pi r^2 dr$$

$$U_{\text{tot}} = \sum_{i < j} u_c(r_{ij}) + \frac{1}{2} N \rho \int_{R_c}^{\infty} u(r) 4\pi r^2 dr$$

$$g(r) \approx 1 \quad \text{for } r > \underline{R_c}$$

$$u_c(r) = \begin{cases} u(r) & r \leq R_c \\ 0 & r > R_c \end{cases}$$



Verlet Integrator

$$\mathbf{r}(t + \Delta t) \approx 2\mathbf{r}(t) - \mathbf{r}(t - \Delta t) + \frac{\mathbf{F}(t)}{m} \Delta t^2$$

- no velocities needed
- same accuracy as velocity Verlet
- velocities can be computed (but $\mathcal{O}(\Delta t^2)$)

$$\mathbf{v}(t) = \frac{\mathbf{r}(t + \Delta t) - \mathbf{r}(t - \Delta t)}{2\Delta t} + \mathcal{O}(\Delta t^2)$$